

UC Irvine

UC Irvine Previously Published Works

Title

Determination of the S-wave π - π amplitude near the peak from the reaction $\pi^- + p \rightarrow \pi^+ + \pi^- + n$

Permalink

<https://escholarship.org/uc/item/3dd9n6p2>

Journal

Physical Review, 155(5)

ISSN

0031-899X

Authors

Bander, M
Shaw, GL

Publication Date

1967-12-01

DOI

10.1103/PhysRev.155.1675

Copyright Information

This work is made available under the terms of a Creative Commons Attribution License, available at <https://creativecommons.org/licenses/by/4.0/>

Peer reviewed

Determination of the S-Wave $\pi\pi$ Amplitude near the ρ Peak from the Reaction $\pi^- + p \rightarrow \pi^+ + \pi^- + n$ *

MYRON BANDER AND GORDON L. SHAW
University of California, Irvine, California

(Received 2 November 1966)

A fit to recent extensive data for the reaction $\pi^- + p \rightarrow \pi^+ + \pi^- + n$ at incident π^- momentum ~ 4 BeV/c and final two-pion center-of-mass energy $m_{\pi\pi} \sim m_\rho$ was made. The peripheral model with absorption was used in the fit. The asymmetry in the final two-pion distribution θ_π gives a quantitative determination of the $\pi\pi$, S-wave, $I=0$ scattering amplitude. A constant phase shift of $\sim +60^\circ$ gives as good a fit to the data as a resonance ϵ^0 (at 730 MeV with a width of 90 MeV), proposed by Durand and Chiu. A negative phase shift of $\sim -60^\circ$ is ruled out by examining the distribution in θ_π as a function of $m_{\pi\pi}$.

I. INTRODUCTION

IT is known¹ that the angular distribution in θ_π for the final two pions in the reaction²

$$\pi^- + p \rightarrow \pi^+ + \pi^- + n \quad (1)$$

near the final two-pion center-of-mass energy $m_{\pi\pi} \sim m_\rho$ requires a large S-wave phase shift³ δ_0 interfering with the $l=1$ production.⁴ Furthermore, the θ_π distribution

of the final pions in the reaction⁵

$$\pi^\pm + p \rightarrow \pi^\pm + \pi^0 + p \quad (2)$$

near $m_{\pi\pi} \sim m_\rho$ yields a small negative value for the $I=2$, S-wave phase shift. Thus, reactions (1) and (2) indicate the presence of a large $\pi\pi$ phase shift δ_0^0 near the ρ region.

The peripheral production model with absorptive corrections gives a good fit⁶⁻⁸ to reaction (2), not only for the cross section as a function of the momentum

* Work supported in part by the National Science Foundation.

¹ See, e.g., G. Shaw and D. Wong, Phys. Rev. **129**, 1379 (1963); M. Islam and R. Piñon, Phys. Rev. Letters **12**, 310 (1964).

² In this paper, we will be discussing data for incident π laboratory momentum ~ 4 BeV/c.

³ A subscript will be used on the amplitudes and phase shifts to denote the l value, and a superscript to denote the isotopic spin.

⁴ At these values for $m_{\pi\pi}$, d waves are neglected (but f^0 production probably becomes important at somewhat higher $m_{\pi\pi}$).

⁵ Saclay-Orsay-Bari-Bologna Collaboration, Nuovo Cimento **25**, 365 (1962).

⁶ K. Gottfried and J. Jackson, Nuovo Cimento **34**, 735 (1964).

⁷ L. Durand and Y. Chiu, Phys. Rev. **137**, B1530 (1965).

⁸ M. Bander and G. Shaw, Phys. Rev. **139**, B956 (1965).

transfer (to the nucleon) t , but to the measured⁹⁻¹¹ angular distributions in θ_π (the polar scattering angle in the two-pion center-of-mass system) and ϕ_π (the Treiman-Yang azimuthal angle). Thus, a similar, detailed calculation of (1) should give quantitative values for δ_0^0 . Such a calculation was performed by Durand and Chiu,¹² who found that the data averaged in $m_{\pi\pi}$ over the ρ peak required an $I=0$, S -wave resonance (ϵ^0) located at 730 MeV with a width of 90 MeV.

More extensive data¹³ are now available, so that we have repeated the calculation of Durand and Chiu in greater detail. We find that not only their resonant solution ϵ^0 , but also a constant δ_0^0 of $\sim \pm 60^\circ$, give equally good fits to the data.^{9-11,13} Furthermore, when the distributions¹³ in θ_π and ϕ_π are fitted as a function of $m_{\pi\pi}$, instead of averaging over the peak, the negative value (or equivalently $\delta_0^0 \sim 120^\circ$) for δ_0^0 is ruled out.¹⁴ (A large negative value for δ_0^0 , as is pointed out by Chew,¹⁵ would have been quite significant with regard to the vacuum Regge trajectory.)

In conclusion, we find that reaction (1) does not require the S -wave resonance ϵ^0 proposed by Durand and Chiu. In addition, the slowly varying value of $\sim +60^\circ$ in the region 650–850 MeV appears to fit smoothly with most of the lower-energy determinations of δ_0^0 .

II. CALCULATIONS AND CONCLUSIONS

We consider the reaction (1) to proceed via the one-pion-exchange diagram shown in Fig. 1 for the production of an S -wave $\pi^+\pi^-$ pair and the neutral ρ . Let the π - π scattering amplitudes which are functions of the invariant $s = m_{\pi\pi}^2$ be $A_0(s)$ and $A_1(s)$. Then the amplitudes for diagrams in Fig. 1 in the peripheral model with absorptive corrections have the form $A_0\langle\lambda|\lambda'\rangle$ and $A_1\langle\lambda|\lambda'\rangle Y_{1^\mu}(\theta_\pi, \phi_\pi)$, where the matrix elements $\langle\lambda|\lambda'\rangle$ and $\langle\lambda|\lambda'\mu\rangle$ are a function of s , the momentum transfer t , and the incident-pion momentum k . Thus, the cross section for (1) is

$$\begin{aligned} \sigma = \sum_{\lambda, \lambda'} \{ & |A_0(s)|^2 |\langle\lambda|\lambda'\rangle|^2 \\ & + 2 \operatorname{Re}[A_0(s)A_1^*(s) \sum_{\mu} \langle\lambda|\lambda'\rangle \langle\lambda|\lambda'\mu\rangle^* Y_{1^\mu}(\theta_\pi, \phi_\pi)] \\ & + \sum_{\mu, \nu} \langle\lambda|\lambda'\mu\rangle \langle\lambda|\lambda'\nu\rangle^* Y_{1^\mu}(\theta_\pi, \phi_\pi) Y_{1^\nu}^*(\theta_\pi, \phi_\pi) |A_1(s)|^2 \} \\ & \times dt ds d\cos\theta_\pi d\phi_\pi. \quad (3) \end{aligned}$$

⁹ Z. Guiragossian, Phys. Rev. Letters **11**, 85 (1963).

¹⁰ V. Hagopian, W. Selove, J. Alitti, J. Baton, and M. Neven-Rene, Phys. Rev. **145**, 1129 (1966).

¹¹ Aachen-Birmingham-Bonn-Hamburg-London (I.C.)-München Collaboration, Phys. Rev. **138**, B897 (1965).

¹² L. Durand and Y. Chiu, Phys. Rev. Letters **14**, 329, 680(E) (1965).

¹³ R. Birge, R. Ely, T. Schumann, Z. Guiragossian, and M. Whitehead, in *Proceedings of the 12th Annual International Conference on High-Energy Physics, Dubna, 1964* (Atomizdat, Moscow, 1965), p. 153; Z. Guiragossian (private communication).

¹⁴ The distribution in θ_π is more sensitive to the S -wave π - π parameters than is the distribution in ϕ_π .

¹⁵ G. Chew, Phys. Rev. **140**, B1427 (1965).

We used the method of Ref. 8 to calculate the amplitudes $\langle\lambda|\lambda'\rangle$ and $\langle\lambda|\lambda'\mu\rangle$. The assumption of total absorption in the relative $l=0$ state of the final two pions and the nucleon was seen to give the best fit to the charged ρ production (2).¹⁶ It is expected that the process in Fig. 1(a) will not be as sensitive to the details of the absorption because of its simpler helicity structure. Thus, we used the same absorption parameters for the diagrams in Fig. 1 as for reaction (2).⁸

The $l=1$, π - π amplitude A_1 was taken as

$$A_1(s) = \frac{1}{\sqrt{2}} \left[\frac{m_\rho^2 - s}{(s - 4m_\pi^2)\gamma_\rho} - i \left(\frac{s - 4m_\pi^2}{s} \right)^{1/2} \right]^{-1}, \quad (4)$$

with $m_\rho = 760$ MeV and γ_ρ corresponding to the width $\Gamma_\rho = 100$ MeV. We included a (small) $I=2$ amplitude as well as the (large) $I=0$ amplitude³ in A_0 :

$$\begin{aligned} (4\pi)^{1/2} A_0 = \frac{1}{\sqrt{3}} & \left[\alpha_0^0(s) - i \left(\frac{s - 4m_\pi^2}{s} \right)^{1/2} \right]^{-1} \\ & + \frac{1}{\sqrt{6}} \left[\alpha_0^2(s) - i \left(\frac{s - 4m_\pi^2}{s} \right)^{1/2} \right]^{-1}, \quad (5) \end{aligned}$$

with

$$\alpha_0^I(s) = \left(\frac{s - 4m_\pi^2}{s} \right)^{1/2} \cot \delta_0^I. \quad (6)$$

$\alpha_0^2(s)$ was taken to be a large negative constant α_0^2 corresponding to $|\delta_0^2| \lesssim 15^\circ$.^{5,17} We considered two forms for δ_0^0 :

$$\alpha_0^0(s) = \alpha_0^0, \quad (7)$$

and

$$\alpha_0^0(s) = (m_0^2 - s)/\gamma. \quad (8)$$

The effects in the present problem of including S -wave π - π production can be observed only in the θ_π and ϕ_π distributions. If these distributions are averaged in $m_{\pi\pi}$ over the ρ peak, we find three types of solutions which give equally good fits to the data: (i) the resonant solution ϵ^0 found by Durand and Chiu; (ii) a large

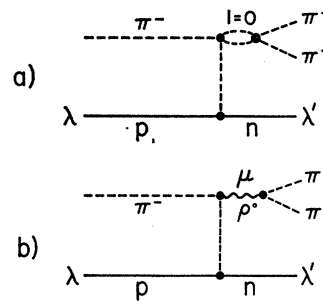


FIG. 1. One-pion exchange diagrams for the reaction (1). λ and λ' are the helicity states of the proton and neutron, respectively. Diagram (a) corresponds to the production of an S -wave π - π pair. Diagram (b) corresponds to the production of the ρ resonance in a helicity state μ .

¹⁶ We note that total absorption in the final ρN state is not necessary if the ρN elastic scattering amplitude is strongly helicity dependent. See M. Bander and G. Shaw, Bull. Am. Phys. Soc. **11**, 23 (1966).

¹⁷ The factor $[(s - 4m_\pi^2)/s]^{1/2}$ is essentially a constant in the energy region we are concerned with, so that $\alpha = \text{constant}$ is approximately the same as $\delta = \text{constant}$.

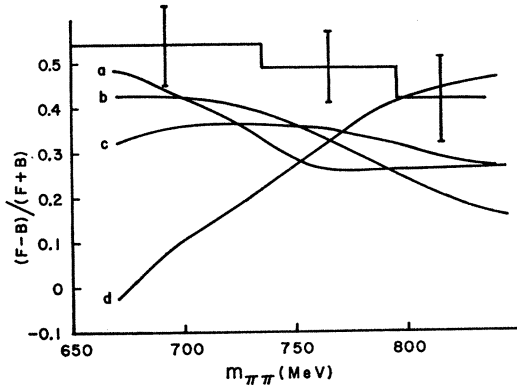


FIG. 2. Plots of the forward-backward asymmetry $(F-B)/(F+B)$ as a function of $m_{\pi\pi}$. The experimental data ($|t| < 10m_{\pi^2}$) are those of Birge *et al.* (Ref. 13). The theoretical curves calculated using the peripheral model with absorption correspond to the S -wave parameters: (a) an $I=0$ resonance at 730 MeV with a width of 100 MeV, no $I=2$ amplitude included; (b) $\alpha_0^0=0.4$ (i.e., $\delta_0^0 \approx 66^\circ$) and no $I=2$; (c) $\alpha_0^0=0.4$ and $\alpha_0^2=-3.0$ (i.e., $\delta_0^2 \approx -17^\circ$); (d) $\alpha_0^0=-0.4$ and no $I=2$.

positive constant δ_0^0 ; (iii) a large negative constant δ_0^0 . Our fits to the data of Birge *et al.*¹³ as a function of $m_{\pi\pi}$ are shown in Figs. 2–4.¹⁸ Although the S -wave contribution is important in fitting the ϕ_π distribution, we note in Fig. 4 that it cannot be used to distinguish

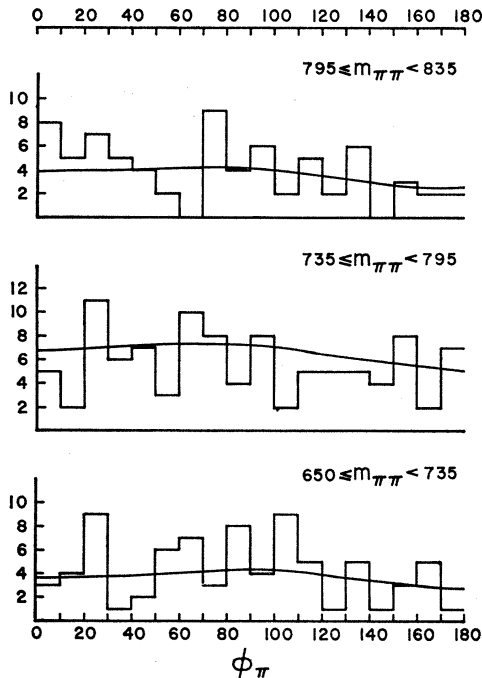


FIG. 3. Plots of the ϕ_π distribution for three bins in $m_{\pi\pi}$. The data are those of Ref. 13. The calculated (smooth) curve corresponds to the S -wave parameters (b) described in Fig. 2.

¹⁸ In addition to the results presented, we calculate the full θ_π, ϕ_π distributions.

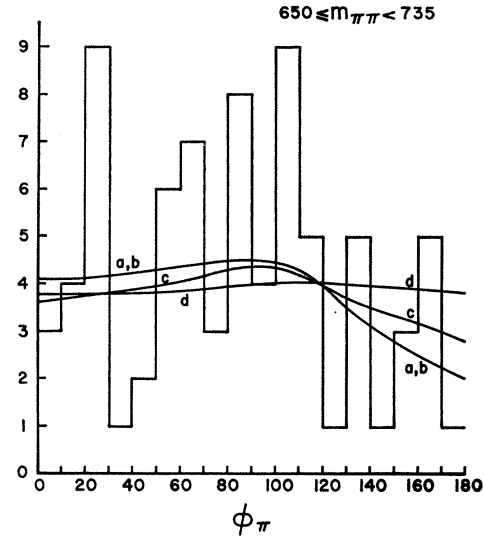


FIG. 4. Plots of ϕ_π in the bin $650 < m_{\pi\pi} < 735$ for the solutions (a)–(d) described in Fig. 2.

between the types of solutions. On the other hand, the forward-backward asymmetry

$$\frac{\sigma(\theta_\pi < \pi/2) - \sigma(\theta_\pi > \pi/2)}{\sigma(\theta_\pi < \pi/2) + \sigma(\theta_\pi > \pi/2)} = \frac{F-B}{F+B}$$

as a function of $m_{\pi\pi}$, shown in Fig. 2, seems to rule out the negative phase-shift solution (d).¹⁹ [Note that as the (negative) $I=2$ phase shift δ_0^2 is made larger in magnitude, the fit with a negative δ_0^0 gets worse.]

Thus, a detailed fit to the data for reaction (1) yields two solutions for δ_0^0 in the energy range $650 \leq m_{\pi\pi} \leq 850$: the narrow resonance ϵ^0 found by Durand and Chiu, and a constant positive value of $\sim 60^\circ$. We feel that the latter, “simpler,” solution is more likely to be correct. This solution seems to fit smoothly with most of determinations of δ_0^0 at smaller $m_{\pi\pi}$.²⁰

ACKNOWLEDGMENTS

We would like to thank Dr. Z. G. T. Guiragossian for helpful discussions and for making the data of Ref. 13 available to us prior to publication. We would like to thank the Physics Department at Stanford University and the Stanford Linear Accelerator Center, where part of this work was done.

¹⁹ Note that *all* the theoretical determinations of $(F-B)/(F+B)$ are lower than the data. However, because of the “background,” the slope of this quantity is probably better determined experimentally than is the absolute normalization.

²⁰ See, e.g., L. Brown and P. Singer, Phys. Rev. **133**, B812 (1964); C. Lovelace, R. Heinz, and A. Donnachie, Phys. Letters **22**, 332 (1966); Y. Fujii, University of Tokyo (unpublished). For a full list of references, see P. Singer, Finnish Summer School, 1966 (to be published).